

# TIME EVOLUTION OF THE REVERSE SHOCK IN SN 1006

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## ABSTRACT

The Schweizer-Middleditch star, located behind the SN 1006 remnant and near its center in projection, provides the opportunity to study cold, expanding ejecta within the SN 1006 shell through UV absorption. Especially notable is an extremely sharp red edge to the Si II 1260 Å feature, which stems from the fastest moving ejecta on the far side of the SN 1006 shell—material that is just encountering the reverse shock. Comparing *HST* far-UV spectra obtained with COS in 2010 and with STIS in 1999, we have measured the change in this feature over the intervening 10.5-year baseline. We find that the sharp red edge of the Si II feature has shifted blueward by  $0.19 \pm 0.05$  Å, which means that the material hitting the reverse shock in 2010 was moving slower by  $44 \pm 11$  km s<sup>−1</sup> than the material that was hitting it in 1999, a change corresponding to  $-4.2 \pm 1.0$  km s<sup>−1</sup> yr<sup>−1</sup>. This is the first observational confirmation of a long-predicted dynamic effect for a reverse shock: that the shock will work its way inward through expanding supernova ejecta and encounter ever slower material as it proceeds. We also find that the column density of *shocked* Si II (material that has passed through the reverse shock) has decreased by  $7 \pm 2\%$  over the ten-year period. The decrease could indicate that in this direction the reverse shock has been ploughing through a dense clump of Si, leading to pressure and density transients.

*Subject headings:* ISM: individual (SN 1006, SNR G327.6+14.6) — shock waves — supernovae: individual (SN 1006) — supernova remnants

## 1. INTRODUCTION

The idea of a reverse shock that can reheat expanding ejecta in a young supernova remnant (SNR) was introduced in a pair of companion papers by Gull (1973) and Rosenberg & Scheuer (1973), and was further articulated by McKee (1974) and by Gull (1975). It was rapidly incorporated into the interpretation of early X-ray data from young SNRs (*e.g.*, Charles et al. 1975), and has become firmly entrenched into the canonical model of a young SNR. Briefly, the idea is that following a supernova (SN) event, a blast wave moves outward through the surrounding interstellar or circumstellar material, which becomes swept up behind the expanding strong shock. Within this shell, a contact discontinuity forms between the expanding SN ejecta and the newly swept-up material, and an internal shock develops that works its way inward through the ejecta, reheating them and giving rise to strong X-ray and optical emission from material that is highly enriched in heavy elements. As the SNR evolves and the reverse shock works its way inward, it

should encounter ever-slower ejecta—a proposition that we test through the observations reported here.

Despite this long pedigree, detailed observations of the development of the reverse shock over time in actual young SNRs have been sparse. The most detailed has been for Cas A: the *Chandra* images clearly showing the reverse shock and its encounter with ejecta of different compositions (Hwang et al. 2000, 2004) and the *Hubble* Space Telescope (*HST*) images showing the evolution of individual ejecta knots as they encounter the reverse shock (Morse et al. 2004; Fesen et al. 2011). In this paper, we report the results of an experiment designed to test a definite prediction of the reverse-shock model: that as the reverse shock moves inward through the freely-expanding ejecta, the velocity of material entering the reverse shock gradually decreases over time. The freely expanding ejecta are too cold to be readily detected in emission, but the velocity profile can be probed through absorption spectroscopy carried out in the UV, where ions commonly found in ejecta have strong resonance lines connecting to the ground state. One of very few objects where this technique has so far been employed is the remnant of the bright Type Ia supernova SN 1006 (SNR G327.6+14.6), located at a distance of 2.2 kpc

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and  $\sim 500$  pc above the Galactic plane (Winkler et al. 2003). Located only  $2'8$  arcmin south of the projected center of the the  $15'$  radius SN 1006 shell, the UV-bright Schweizer & Middleditch (1980, henceforth SM) star enables study of a “core sample” through the remnant.<sup>1</sup>

The SM star is an OB subdwarf, probably not much farther away than SN 1006. It is relatively bright ( $V = 16.74$ ,  $B - V = -0.14$ ), with low foreground extinction,  $E(B - V) = 0.1$  (Schweizer & Middleditch 1980; Wu et al. 1993). Its potential as a background “UV light-bulb” for probing the interior of SN 1006 through absorption spectroscopy was first exploited by Wu et al. (1983) using the International Ultraviolet Explorer (*IUE*), with subsequent studies carried out by Fesen et al. (1988); Wu et al. (1993, 1997), and Blair et al. (1996). These spectra show strong Fe II lines at  $2383 \text{ \AA}$  and  $2600 \text{ \AA}$  that are broad and nearly symmetric, and in addition lines from Si II, Si III, and Si IV, all of which are also broad, but with profiles that are entirely at positive (redshifted) velocities (Wu et al. 1997). The most detailed UV spectrum of the SM star was obtained from the *Hubble* Space Telescope (*HST*) in 1999 with the Space Telescope Imaging Spectrograph (STIS) and reported by Hamilton et al. (2007, hereafter HFB07), which also summarizes the previous UV studies.

The strongest feature in the  $1150\text{--}1700 \text{ \AA}$  band is broad absorption due to Si II  $1260.4 \text{ \AA}$ , all on the red side of the line, extending from a velocity of  $\sim +2400 \text{ km s}^{-1}$  to a sharp edge at  $+7026 \text{ km s}^{-1}$  (HFB07). Hamilton et al. (1997, 2007) have modeled this profile as a combination of both shocked and unshocked Si II ejecta within SN 1006: the highest redshifts result from ejecta that are still freely expanding on the far side of the shell, with the sharp edge due to material just before it enters the reverse shock. The reverse-shock model predicts that, over time, the shock encounters ever slower-moving ejecta, and hence the red edge should gradually move to shorter wavelength. We report here a second-epoch spectrum of the SM star obtained with the Cosmic Origins Spectrograph (COS) on *HST* that, when compared with the 1999 STIS spectrum, demonstrates this phenomenon for the first time.

## 2. OBSERVATIONS AND WAVELENGTH REGISTRATION

For this study, we used as the first-epoch spectrum the far-ultraviolet echelle spectrum of the SM star obtained with STIS in July 1999 and described by HFB07. Briefly, the STIS observation used a  $2'' \times 2''$  slit, the E140M grating and the FUV-MAMA detector to cover wavelengths  $1150\text{--}1700 \text{ \AA}$  at a scale of  $0.0122 \text{ \AA pixel}^{-1}$ . The line spread function is non-Gaussian, with a narrow core and broad wings, and in the crucial region of the spectrum near  $1300 \text{ \AA}$  has a FWHM of  $4.6 \text{ km s}^{-1}$ .

For the second-epoch observation, we used the newly installed COS instrument, since it has a far higher throughput than STIS, though with spectral resolution that is almost four times lower. The observation was carried out on 2010 Jan 16 using the  $2'5$  diameter primary science aperture and the G130M grating. A central wavelength setting of  $1327 \text{ \AA}$  and standard subsampling

(FP-POS = auto) were used to cover the range  $\sim 1170\text{--}1470 \text{ \AA}$ , with a gap extending over  $1317\text{--}1327 \text{ \AA}$ , at a dispersion of  $0.010 \text{ \AA pixel}^{-1}$ . The COS line spread function is also non-Gaussian, with FWHM  $18 \text{ km s}^{-1}$  near  $1300 \text{ \AA}$  (see Sahnou et al. 2010, for a summary of the COS on-orbit performance). Because of the higher throughput with COS, only a single orbit was required to obtain a spectrum with signal-to-noise similar to that obtained in eight orbits with STIS. Key aspects of the two observations are summarized in Table 1, and in Fig. 1 we plot the spectra from both observations, binned to emphasize the broad features. The time separation between the epochs is  $10.48 \text{ yr}$ .

In order to measure the small shift in the wavelength of the absorption edge, it is crucial that the STIS and COS spectra be precisely registered in wavelength. In order to achieve more precise registration than can be done using the wavelength scale provided with pipeline-processed COS data, we can use the many narrow interstellar and stellar lines in the wavelength region near the edge. We have used a cross-correlation analysis (the IRAF<sup>2</sup> task `fxcor`), first to cross-correlate each of the four individual (subsampling) frames, made tiny wavelength shifts to best align these, and averaged them to optimize the total COS spectrum. We then used a similar cross-correlation analysis to determine the small offsets between the COS and STIS spectra on both the blue ( $1245\text{--}1264 \text{ \AA}$ ) and red ( $1292\text{--}1301.5 \text{ \AA}$ ) sides of the edge, and found that with the original wavelength scales, the narrow lines differ in wavelength by as much as  $40 \text{ m\AA}$ , or  $9.6 \text{ km s}^{-1}$ . The differences are systematic, so we were able to apply a simple linear transformation to the COS wavelength scale to achieve registration with the STIS one to within  $\pm 6 \text{ m\AA}$ , or  $\pm 1.5 \text{ km s}^{-1}$ . In Fig. 2 we plot the difference between measured wavelengths in the COS and STIS spectra for several narrow lines, both before and after the transforming the COS spectrum. After registering the COS and STIS spectra, we find a clear shift in the red edge of the  $1260 \text{ \AA}$  absorption feature to shorter wavelength, which we have measured as described in the next section.

## 3. MODEL FIT TO THE COS SPECTRUM

As discussed by Hamilton et al. (1997, 2007), the unusual shape of the broad, redshifted Si II  $1260 \text{ \AA}$  absorption feature in SN 1006 can be understood as arising from a combination of unshocked and shocked Si II ejecta. The sharp red edge at  $\sim 7000 \text{ km s}^{-1}$  marks the position of a reverse shock front, where unshocked Si II is “instantaneously” decelerated and thermalized. The lifetime of the resulting shocked Si II against collisional ionization is of the order of the age of the remnant, so it is not surprising that shocked Si II should survive to remain observable. While separated three-dimensionally of course, both shocked and unshocked Si II lie along the line of sight.

Fig. 3 shows the result of fitting the  $1260 \text{ \AA}$  feature observed with COS to a combination of unshocked and shocked Si II. The analysis is essentially the same as

<sup>1</sup> In addition to the SM star, Winkler et al. (2005) have used two fainter background QSOs to probe additional lines of sight through SN 1006.

<sup>2</sup> IRAF is distributed by the National Optical Astronomy Observatories, which is operated by the AURA, Inc. under cooperative agreement with the National Science Foundation.

that of HFB07, to which the reader is referred for more detail. Fig. 3 can be compared to Figure 4 of HFB07, which showed a similar fit to the same feature observed with STIS.

In brief, each of the observed spectra were interpolated across narrow interstellar lines, and their continua fitted to a low-order polynomial. In the case of the STIS spectra of the SM star and the comparison star PG 0839+399 analysed in 2007, the continuum was a quintic polynomial in log flux versus inverse wavelength fitted over unabsorbed regions from 1150 to 1700 Å. The COS and STIS spectra are in generally excellent agreement (see §5), but there is a small but significant difference in the overall level of fluxes. To compensate for this difference, we multiplied the COS spectrum by a factor linear in log flux versus inverse wavelength, varying from 1.03 at 1190 Å to 1.08 at 1290 Å.

Fig. 3 also shows the ratio of the interstellar-line-excised, continuum-corrected spectrum of the SM star observed with COS, to that of the comparison star PG 0839+399 observed with STIS. The comparison star PG 0839+399 (Green, Schmidt, & Liebert 1986) was identified by Blair et al. (1996) as having not only a similar temperature and gravity, but also similar photospheric abundances and low extinction. As remarked by HFB07, the spectrum of the comparison star is not as similar in detail to that of the SM star as one might have liked, but dividing the SM spectrum by the comparison one does help to clarify which features in the SM star spectrum are due to SN 1006.

Three separate velocity parameters can be measured from the line profile shown in Fig. 3: the free-expansion velocity  $r_s/t$  (shock radius  $r_s$  divided by age  $t$ ) of unshocked Si II at the reverse shock front, and the mean velocity  $\bar{v}$  and one-dimensional velocity dispersion  $\sigma$  of shocked Si II. Not surprisingly, in view of the good agreement between the STIS and COS spectra, and the similar analysis, the parameters measured from COS are quite similar to those measured from STIS, Table 2.

The three velocities are related by shock jump conditions. As noted by HFB07, the measured velocities point to essentially all of the shock energy going into heating ions, with little energy in heating electrons, consistent with other observational evidence (Laming et al. 1996; Ghavamian et al. 2002; Vink et al. 2003; Vink 2005), and also with little energy in accelerating particles (e.g. Kosenko et al. 2011). In the absence of electron heating or particle acceleration, the change  $\Delta v$  in the velocity between the unshocked and shocked Si II should equal the three-dimensional velocity dispersion  $\sqrt{3}\sigma$  of the shocked Si II,

$$\Delta v \equiv \frac{r_s}{t} - \bar{v} = \sqrt{3}\sigma. \quad (1)$$

If some of the shock energy went into electron heating or particle acceleration, then the velocity dispersion would be correspondingly smaller. If all three velocities are treated as free parameters, the observations mildly prefer the opposite, an ion velocity dispersion slightly ( $1\sigma$ ) larger than predicted by equation (1). In the present paper we choose to constrain the fit, Fig. 3, such that equation (1) is satisfied, whereas HFB07 chose to measure all three velocities separately from the STIS spectrum. HFB07 estimated a  $3\sigma$  upper limit of 0.26 of the

shock energy going into forms other than ion heating.

Fits constrained to satisfy equation (1) prefer slightly less shocked Si II absorption at its peak near  $5,000 \text{ km s}^{-1}$ . A simple way to accommodate the preference is to fit the blue edge of the profile up to  $5,100 \text{ km s}^{-1}$  instead of the  $5,600 \text{ km s}^{-1}$  adopted by HFB07. The values quoted in Table 2 are best fits over the reduced range. The quoted uncertainties are  $1\sigma$  statistical uncertainties about the best fit, subject to variations *not* constrained by equation (1).

The observed column density of shocked Si II measured from the line profile in Fig. 3 is (compare eq. 8 of Hamilton et al. 1997)

$$N_{\text{SiII}}^{\text{shk}} = \frac{m_e c}{\pi e^2 f \lambda} \sqrt{2\pi} \sigma \tau_0 = 7.4 \pm 0.3 \times 10^{14} \text{ cm}^{-2} \text{ (measured),} \quad (2)$$

where  $\lambda = 1260 \text{ Å}$  is the wavelength of the line,  $f = 1.007$  (Morton 1991) is its oscillator strength,  $\tau_0$  is the optical depth at line center of the fitted Gaussian profile of shocked Si II, and the factor  $\sqrt{2\pi}$  comes from integrating over the profile. The uncertainty in the estimate (2) is the statistical uncertainty arising from photon shot noise, and does not include “systematic” uncertainty arising from modeling choices, which include placement of the continuum, and the range of velocities over which the fitting is done. The systematic uncertainty is somewhat larger. Notably, if we extend the range over which we fit the blue edge of the shocked Si II profile to  $5,600 \text{ km s}^{-1}$  (as in HFB07) instead of  $5,100 \text{ km s}^{-1}$ , then the measured column density increases to  $8.0 \times 10^{14} \text{ cm}^{-2}$ . The measured value (2) is listed in Table 2, along with the value that HFB07 would have measured from STIS but did not report.

The measured column density (2) of shocked Si II is consistent with that predicted by steady ionization of Si II (compare eq. 13 of Hamilton et al. 1997),

$$N_{\text{SiII}}^{\text{shk}} = \frac{n_{\text{SiII}} v_s}{4n_e \langle \sigma v \rangle_{\text{SiII}}} \\ = 6 \pm 2 \times 10^{14} \text{ cm}^{-2} \left( \frac{n_e/n_{\text{SiII}}}{2.5} \right)^{-1} \text{ (predicted),} \quad (3)$$

where  $\langle \sigma v \rangle_{\text{SiII}}$  is the collisional ionization rate (see below), and the quoted uncertainty is a nominal estimate of the uncertainty in that rate. The electron to Si II ratio of  $n_e/n_{\text{SiII}} = 1.5/0.6 = 2.5$  in equation (3) is a fiducial value that assumes that each Si ion contributes on average 1.5 electrons during its steady state ionization, and that the Si is in a silicon-rich region where it comprises 60% of the ion abundance by number, the remaining 40% being mostly S, consistent with SN Ia nucleosynthesis models (Höflich et al. 1998).<sup>3</sup>

As discussed by Hamilton et al. (1997), it is easy to reduce the theoretical value (3) of shocked Si II, by diluting the Si with other elements (i.e. by increasing  $n_e/n_{\text{SiII}}$ ), or by truncating the Si downstream, but it is harder to increase the theoretical value, to approach the slightly greater observed value (2). Thus agreement between the measured and predicted column densities of shocked Si II appears to require that the absorbing shocked Si

<sup>3</sup> Many lines of evidence suggest that SN 1006 was a Type Ia event (e.g., Winkler et al. 2003).

arises from silicon-rich ejecta. This conclusion is consistent with the absence of Fe II absorption with the same profile as the shocked Si II. Although Fe II absorption is observed, its profile is dramatically different from that of the shocked Si II, indicating that little iron is mixed in with the shocked silicon (Hamilton et al. 1997).

In Hamilton et al. (1997) there seemed to be a possible discrepancy between the observed and predicted values of the column density of shocked Si II. That discrepancy has now disappeared, in part thanks to a reduction in the collisional ionization cross-section. The collisional ionization rate used in equation (3) is  $\langle\sigma v\rangle_{\text{SiII}} = 4.4 \times 10^{-8} \text{ cm}^3 \text{ s}^{-1}$ , obtained by integrating cross-sections from Clark & Abdallah (2003) over a Maxwellian distribution of colliding electrons at temperature 95 eV, the temperature reached by electrons as a result of Coulomb collisions with Si II ions over the collisional ionization timescale of Si II. The collisional ionization rate from Clark & Abdallah (2003) is smaller than the rate  $6.1 \times 10^{-8} \text{ cm}^3 \text{ s}^{-1}$  from Lennon et al. (1988) used by Hamilton et al. (1997). The newer rate, based on theoretical computations, should be more reliable than the older rate, based on isoelectronic scaling.

A concern is that the collisional ionization time  $t_{\text{SiII}}$  of Si II is comparable to the age  $t$  of the remnant (specifically,  $t_{\text{SiII}}/t \approx 0.8$ , given the pre-shock Si II density from the observed optical depth of freely expanding unshocked Si II at the shock front, an estimated electron-to-Si II density  $n_e/n_{\text{SiII}} \approx 2.5$ , and the theoretical collisional ionization rate (cf. Hamilton et al. 1997)). The collisional ionization time is long enough that the shocked Si II absorption should be a superposition of components shocked at various times with various velocities. It is somewhat surprising that the profile should fit well to a single component satisfying the shock jump conditions. A possible explanation of the coincidence is that one-dimensional hydrodynamic models predict rather flat velocity and temperature profiles of shocked ejecta (e.g. Dwarkadas & Chevalier 1998; Kosenko et al. 2011). Another possible explanation, suggested by the faster than expected change in the optical depth of shocked Si II, §5, is that the absorption is being produced by a lump of silicon-rich ejecta whose average density was higher than that observed today, and which has been shocked over a timescale shorter than the age of the remnant.

#### 4. EVOLUTION OF THE SHOCK FRONT FROM STIS TO COS

As the reverse shock propagates into the Si II ejecta, the free-expansion velocity  $r_s/t$  of unshocked gas entering the shock at radius  $r_s$  should decrease with time  $t$  as

$$\frac{d(r_s/t)}{dt} = -\frac{v_s}{t}, \quad (4)$$

where  $v_s$  is the velocity of the reverse shock relative to the ejecta. The shock velocity  $v_s$  is related to the change  $\Delta v$  in the velocity between unshocked and shocked Si II by  $v_s = (4/3)\Delta v$  (regardless of electron heating). The deceleration  $\Delta v$  measured from the COS spectrum shown in Fig. 3 yields a shock velocity of  $v_s = 2630 \pm 120 \text{ km s}^{-1}$ , in good agreement with the STIS measurement of  $v_s = 2680 \pm 120 \text{ km s}^{-1}$ . As described by HFB07, equation (4) coupled with the measured shock velocity predicts that the free-expansion velocity at the shock should be de-

creasing at the rate of  $-2.7 \pm 0.1 \text{ km s}^{-1} \text{ yr}^{-1}$ . The predicted change  $\Delta(r_s/t)$  in the free-expansion velocity of the shock over the 10.5 year interval between the STIS and COS observations is then

$$\Delta(r_s/t) = -28 \pm 1 \text{ km s}^{-1} \quad (\text{predicted}). \quad (5)$$

Fig. 4 compares the COS and STIS spectra in the vicinity of the putative reverse shock front near  $7000 \text{ km s}^{-1}$ . To permit direct comparison of the spectra, the COS spectrum has been convolved with the STIS line spread function (LSF), while the STIS spectrum has been convolved with the COS LSF. The net resolution is essentially that of COS, since the COS resolution ( $18 \text{ km s}^{-1}$  FWHM) is over three times broader than STIS ( $4.6 \text{ km s}^{-1}$  FWHM). The flux in the COS spectrum has been multiplied by a slowly-varying factor, about equal to 1.08 near  $1290 \text{ \AA}$ , to bring it to the same normalization as STIS, as described in §3.

The spectra in Fig. 4 show that the shock front near  $7000 \text{ km s}^{-1}$  has indeed moved during the interval between the STIS and COS observations. One way to measure the velocity change would be to shift the two spectra horizontally, but the plethora of stellar lines precludes a precise measurement by this naive approach.

We follow instead a more rigorous approach that models the physics of the situation and that is unbiased by the presence of stellar lines. Fig. 3 shows that, at the shock front, unshocked gas is attenuating the stellar flux by a factor of 0.4. If this is correct, and the shock front has shifted by  $\Delta(r_s/t)$  from STIS to COS, then the STIS spectrum should equal the COS spectrum multiplied by an inverted top hat  $1 - 0.4 = 0.6$  deep and  $\Delta(r_s/t)$  wide. We can think of the top hat as representing material that was absorbing in 1999 that has been “eaten away” by the reverse shock since. Observationally, this absorbing top hat will be convolved with the overall LSF of the spectrum.

Fig. 4 shows the best-fit COS spectrum adjusted by this procedure, with the convolved, absorbing top hat shown as a dotted line. The fit was obtained by minimizing the  $\chi^2$  between the STIS and adjusted COS spectra, the noise being taken to be photon shot noise. The correlations between spectral bins introduced by convolving the spectra were taken into account.

The best fitting velocity shift between STIS (1999) and COS (2010) spectra is

$$\begin{aligned} \Delta(r_s/t) &= -44 \pm 11 \text{ km s}^{-1}, \\ \text{or } &-4.2 \pm 1.0 \text{ km s}^{-1} \text{ yr}^{-1} \quad (\text{measured}). \end{aligned} \quad (6)$$

The measured  $1\sigma$  uncertainty of  $11 \text{ km s}^{-1}$  is about 60% of a COS resolution element. There can be little doubt that the velocity shift is real, as the measured value differs from zero by  $4\sigma$ . The measured velocity shift (6) is reasonably consistent with the predicted shift (5), differing by  $16 \pm 11 \text{ km s}^{-1}$ , or  $1.5\sigma$ . The value of  $6982 \text{ km s}^{-1}$  for  $r_s/t$  given in Table 2 is just the STIS value (which is more precisely measured) reduced by  $44 \text{ km s}^{-1}$ .

It should be remarked that we have assumed in this paper, as previously in Hamilton et al. (1997, 2007), that the shock front is perpendicular to the line of sight, a plausible assumption given that the SM star is close to the projected center of the SN 1006 remnant. If the shock

front normal were inclined by an angle  $\theta$  to the line of sight, but the ejecta are moving directly away from us, then the main effect would be to increase the observed rate of change of the free-expansion velocity by a factor  $1/\cos\theta$ . If a future observation should measure a change faster than that predicted by equation (5) by a statistically significant margin, then one natural explanation would be that the shock front is not perpendicular to the line of sight.

#### 5. APPARENT DECREASE IN SHOCKED Si II ABSORPTION

An unanticipated change between the two spectra is that the optical depth in shocked Si II appears to have decreased over the 10.5 year interval between STIS and COS observations. The change is illustrated in Fig. 5, which compares the COS (2010) and STIS (1999) spectra over the full range of the broad, redshifted Si II absorption feature. The ratio of COS to STIS fluxes indicates a small but significant flux excess, with a profile similar to that of the shocked Si II illustrated in Fig. 3. As discussed in §3, the COS flux has first been multiplied by a linear function varying from 1.03 at 1190 Å to 1.08 at 1290 Å in order to bring its *overall* flux to the same level as STIS.

To measure the apparent change in absorption, we first adjusted the COS spectrum for the  $44 \text{ km s}^{-1}$  shift in the shock front as described in §4, so that the two should be comparable. We then assumed that the COS:STIS ratio has the same profile as shocked Si II but with unknown amplitude, and did a one-parameter least-squares fit to determine that amplitude. This procedure is more precise than simply taking the ratio of shocked column densities in Table 2, whose values are model-dependent. We find that the column density of shocked Si II measured from the COS spectrum is  $0.93 \pm 0.02$  times that in the STIS spectrum from ten years earlier. The statistically significant  $7 \pm 2\%$  change is larger than the  $\sim 1\%$  change one might simplistically have expected over 10 years of observation of a 1000 year old supernova remnant. The 2% uncertainty includes only statistical uncertainty, and not any systematic uncertainty associated with our modeling of the continuum. The 7% change is comparable to the  $\sim 8\%$  systematic adjustment we have made to the COS flux, and some systematic uncertainty attends our modeling of the relative sensitivities between COS and STIS.

If the unexpectedly rapid decrease in optical depth of shocked Si II is real, then a possible explanation is that the shock has been ploughing through a clump of Si. Numerical simulations of shocks overrunning inter-

stellar clouds (Klein et al. 1994) indicate that clouds of sufficiently high density contrast can undergo significant pressure and density transients as the shock propagates through the cloud. The line profile of *unshocked* Si II illustrated in Figure 3 does indicate a density decreasing inward by a factor of 2 over an interval of about  $700 \text{ km s}^{-1}$ . This change is equivalent to about 2% in  $30 \text{ km s}^{-1}$ , which is several times smaller than required by the observed change in optical depth. The profile of unshocked Si II shown in Figure 3 does show possible smaller scale structure consistent with more rapid variation, but the reality of such structure is confused by the many stellar lines.

#### 6. SUMMARY

Using UV absorption-line spectra against the SM star taken with STIS in 1999 and with COS in 2010, we have probed changes in the absorbing supernova ejecta seen along a “core sample” through the SN 1006 remnant. At both epochs we see strong absorption due to Si II, all of it red-shifted and with a sharp red edge. Following Hamilton et al. (1997, 2007), we have interpreted the profile of the strong 1260 Å absorption feature as resulting from both unshocked and shocked (by the reverse shock) Si II ejecta, with the sharp red edge to this feature being due to freely expanding ejecta just before being hit by the reverse shock. After carefully taking into account small systematic differences in both wavelength and flux between measurements with the two instruments, we find that the red absorption edge has shifted by  $-44 \pm 11 \text{ km s}^{-1}$  over the 10.5 yr interval between observations, or  $-4.2 \pm 1.0 \text{ km s}^{-1} \text{ yr}^{-1}$ . This measurement confirms the long-held expectation that in a young SNR, the reverse shock will encounter ever-slower material as it works its way inward through the supernova ejecta. We also find the unanticipated result that the optical depth of shocked Si II has decreased by  $7 \pm 2\%$ , faster than might be expected to result from homologous expansion over 10 years in the SNR’s 1000-yr history. Clumped Si ejecta provides a possible explanation both for the decrease in optical depth and for the absence of corresponding blue-shifted Si II absorption.

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TABLE 1  
OBSERVATIONS OF THE SCHWEIZER-MIDDLEDITCH STAR FROM *HST*

Instrument	Obs. Date	Grating	FWHM ( $\text{km s}^{-1}$ )	Exposure (s)	<i>HST</i> Program ID
STIS	1999 Jul 24-25	E140M	4.6	23,186	7349
COS	2010 Jan 16	G130M	18.	1,756	11659

TABLE 2  
MEASURED PARAMETERS

Parameter	COS (this paper)	STIS (HFB07)
$r_s/t$	$6982 \pm 11 \text{ km s}^{-1}$	$7026 \pm 3 \text{ (rel.)} \pm 10 \text{ (abs.) km s}^{-1}$
$\bar{v}$	$5030 \pm 70 \text{ km s}^{-1}$	$5160 \pm 70 \text{ km s}^{-1}$
$\sigma$	$1140 \pm 50 \text{ km s}^{-1}$	$1160 \pm 50 \text{ km s}^{-1}$
$v_s$	$2630 \pm 120 \text{ km s}^{-1}$	$2680 \pm 120 \text{ km s}^{-1}$
$N_{\text{SiII}}^{\text{shk}}$	$7.4 \pm 0.3 \times 10^{14} \text{ cm}^{-2}$	$8.6 \pm 0.3 \times 10^{14} \text{ cm}^{-2}$

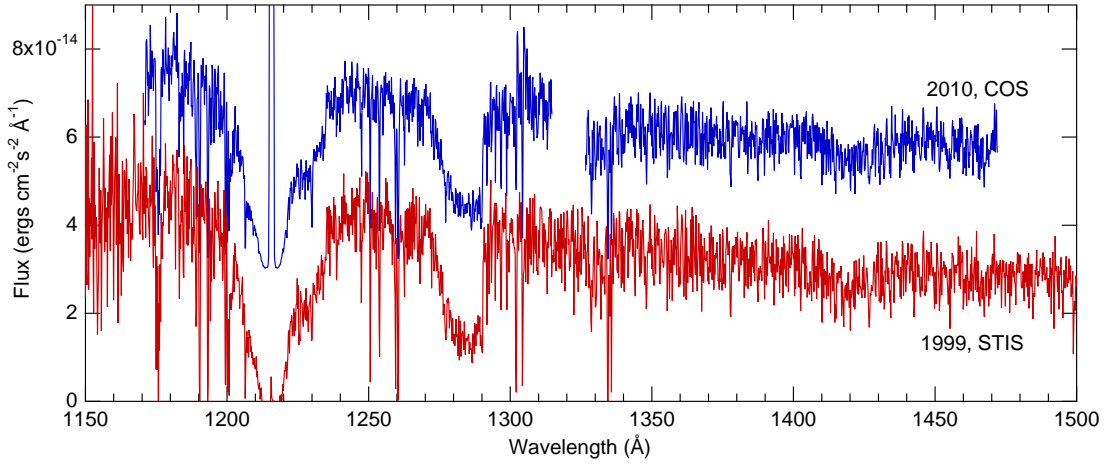


FIG. 1.— The complete COS and STIS FUV spectra of the SM star. The absorption feature at 1217 Å has contributions from geocoronal Ly $\alpha$  emission, stellar and interstellar Ly $\alpha$  absorption, redshifted Si II 1190.4 Å and 1193.3 Å absorption, and redshifted Si III 1206.5 Å absorption. The most prominent “clean” feature is the broad, redshifted Si II 1260.4 Å line, due to a combination of shocked and unshocked Si II ejecta. The COS spectrum has been displaced vertically for clarity.

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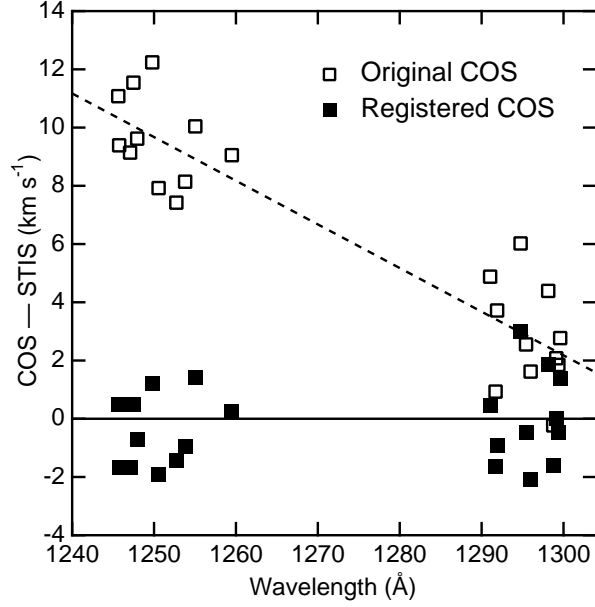


FIG. 2.— The difference between measured wavelengths of several narrow lines, on both blue and red sides of the Si II absorption feature, is plotted as a function of wavelength. A simple linear transformation in the COS wavelength scale brings the two spectra into registration, with an RMS uncertainty of  $\pm 1.5 \text{ km s}^{-1}$ .

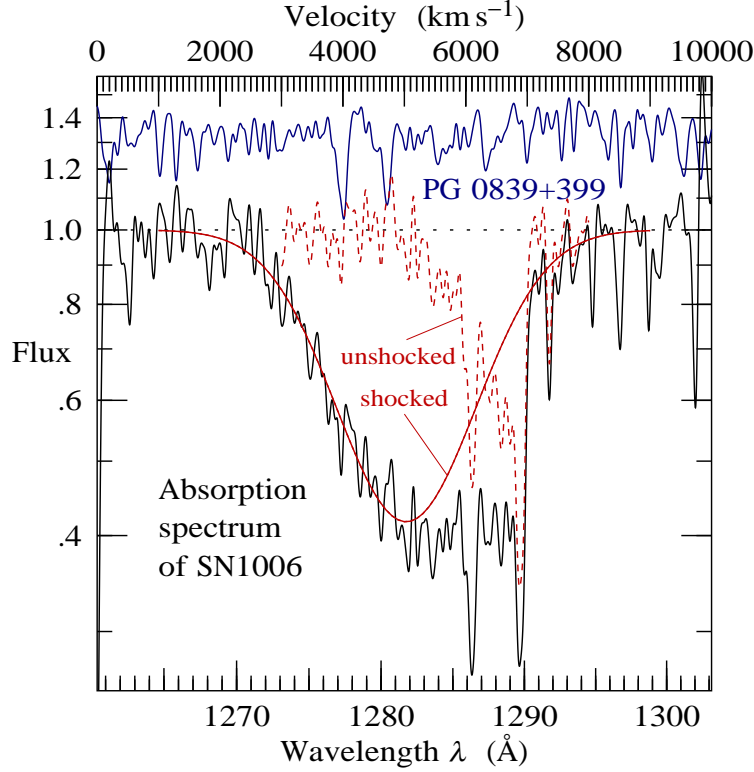


FIG. 3.— COS absorption spectrum of SN 1006 around the redshifted Si II 1260 Å feature, showing the best-fit Gaussian profile of shocked Si II, and the residual unshocked Si II. The spectrum is the ratio of the interstellar-line-excised, continuum-corrected COS spectrum of the SM star to the STIS spectrum of the comparison star PG 0839+399, with each smoothed to a resolution of  $80 \text{ km s}^{-1}$  FWHM before their ratio was taken. The upper spectrum shows the interstellar-line-excised, continuum-corrected comparison stellar spectrum of PG 0839+399 at the same resolution, shifted by  $+32 \text{ km s}^{-1}$  to mesh its stellar lines with those of the SM star, and offset vertically to separate it from the SN 1006 spectrum.

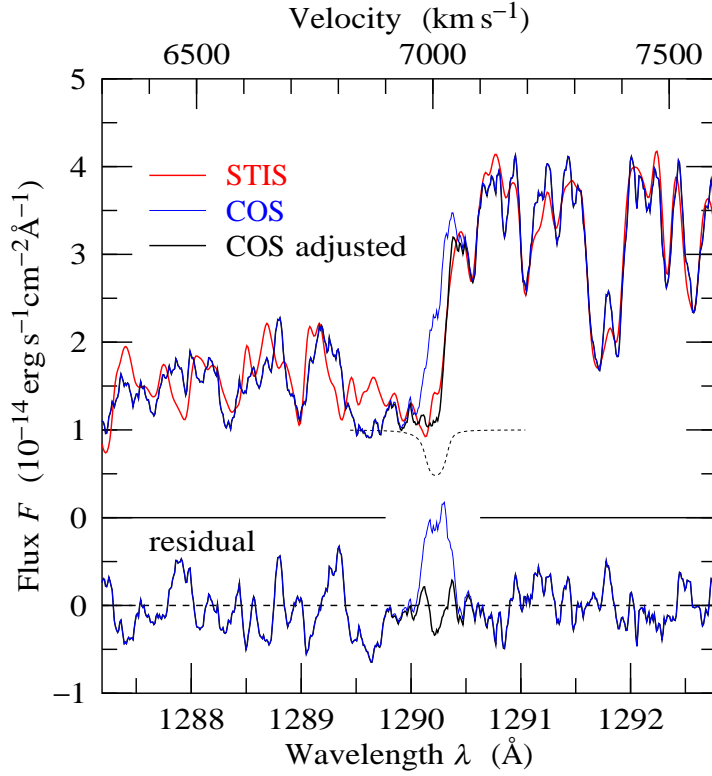


FIG. 4.— Comparison of COS (blue) and STIS (red) spectra of the SM star behind SN 1006. To permit direct comparison, each spectrum has been convolved with the line spread function of the other. Also shown is the best-fit adjusted COS spectrum (black), in which the COS flux has been attenuated by a factor of 0.4 over a slice of velocity  $44 \text{ km s}^{-1}$  wide starting at the  $7026 \text{ km s}^{-1}$  velocity of the shock front observed with STIS. The dotted line shows the factor by which the COS spectrum has been multiplied, an inverted top hat convolved with the overall LSF. The residual difference between the COS and STIS spectra is shown.

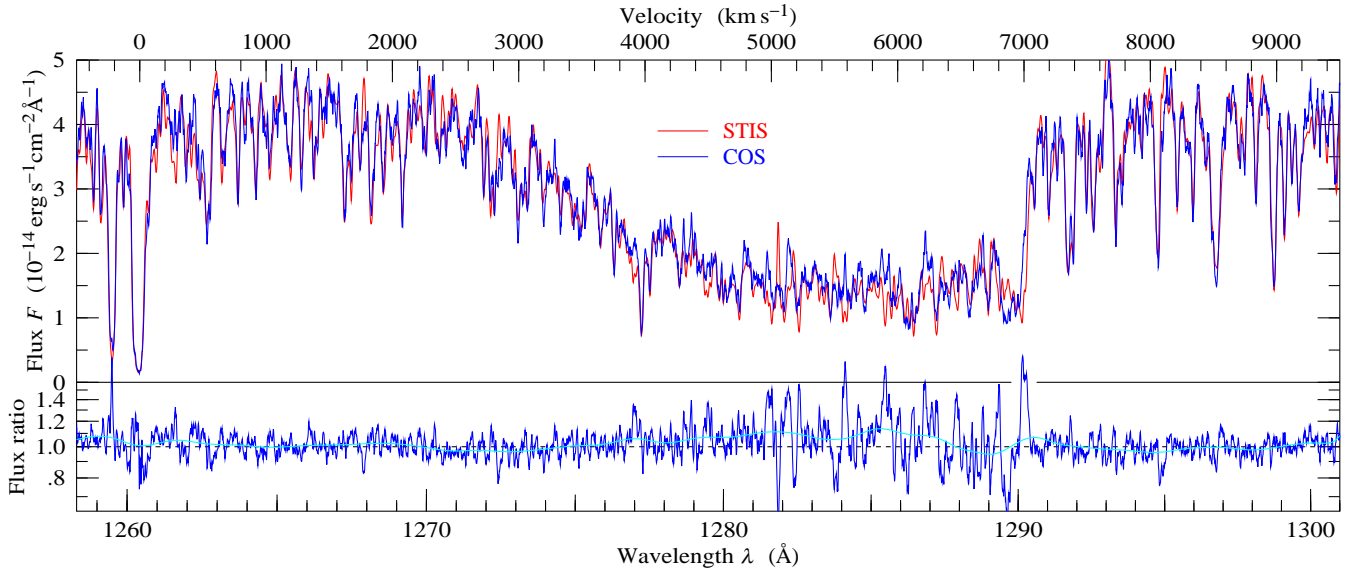


FIG. 5.— Comparison of COS (blue) and STIS (red) spectra of the SM star behind SN1006. Each spectrum has been convolved with the line spread function of the other. The lower panel shows the ratio of the COS to STIS fluxes. The smooth (cyan) line is the ratio smoothed with a Gaussian to  $320 \text{ km s}^{-1}$  FWHM. The ratio shows a small but significant excess with a profile similar to that of the shocked Si II illustrated in Figure 3, indicating a decrease in the optical depth of shocked Si II in COS compared to STIS.